

## In situ and ground-based intercalibration measurements of plasma density at $L = 2.5$

M. A. Clilverd,<sup>1</sup> F. W. Menk,<sup>2</sup> G. Milinevski,<sup>3</sup> B. R. Sandel,<sup>4</sup> J. Goldstein,<sup>5</sup>  
 B. W. Reinisch,<sup>6</sup> C. R. Wilford,<sup>7</sup> M. C. Rose,<sup>1</sup> N. R. Thomson,<sup>8</sup> K. H. Yearby,<sup>9</sup>  
 G. J. Bailey,<sup>7</sup> I. R. Mann,<sup>10</sup> and D. L. Carpenter<sup>11</sup>

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[1] Two independent ground-based experiments and two satellite-borne experiments are used to interpret the changes in plasmaspheric composition at the same point in space during moderate geomagnetic activity on 22 January and 14 February 2001. Mass density at  $L = 2.5$  was determined from an array of magnetometers on the Antarctic Peninsula, while the electron number density along the same flux tube was determined from analysis of the group delay of man-made VLF transmissions from north-east America. The IMAGE satellite RPI experiment provided in situ measurements of the electron number density in passing the equatorial region of the same field line, while the EUV Imager experiment was able to resolve the  $\text{He}^+$  abundance by looking back toward the same place a few hours later. On 22 January 2001 all measurements were consistent with a moderately disturbed plasmasphere. On 14 February 2001 there appeared to be a significant response of the plasmasphere to the moderate ( $K_p = 5$ ) activity levels. Both the electron number density and the mass density determined from the ground-based experiments were markedly higher than on 22 January 2001. Also, the IMAGE RPI gave a markedly lower electron number density than did the ground-based data; this is explained by differences in the longitude at which the measurements were made and the presence of localized plasmaspheric structures. At Antarctic Peninsula longitudes a  $\text{He}^+$  column abundance value of  $6 \times 10^{10} \text{ cm}^{-2}$  is found to be equivalent to plasmaspheric electron density levels of  $3000 \text{ cm}^{-3}$  at  $L = 2.5$ . For these conditions the  $\text{He}^+$  mass abundance was about 12–16% compared with  $\text{H}^+$ . Both decreases and increases in the  $\text{He}^+$  column abundance measured by the EUV Imager appear to be linearly correlated to changes in the percentage occurrence of  $\text{He}^+$  as determined from a combination of ground-based VLF and ULF observations. *INDEX TERMS:* 2768

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<sup>1</sup>British Antarctic Survey, Cambridge, UK.

<sup>2</sup>School of Mathematical and Physical Sciences and Cooperative Research Centre for Satellite Systems, University of Newcastle, Callaghan, Australia.

<sup>3</sup>Ukrainian Antarctic Center, Kyiv, Ukraine.

<sup>4</sup>Lunar and Planetary Laboratory, University of Arizona, Tucson, Arizona, USA.

<sup>5</sup>Rice Space Institute, Rice University, Houston, Texas, USA.

<sup>6</sup>Center for Atmospheric Research, University of Massachusetts, Lowell, Lowell, Massachusetts, USA.

<sup>7</sup>Space and Atmosphere Research Group, Department of Applied Mathematics, University of Sheffield, Sheffield, UK.

<sup>8</sup>Department of Physics, University of Otago, Dunedin, New Zealand.

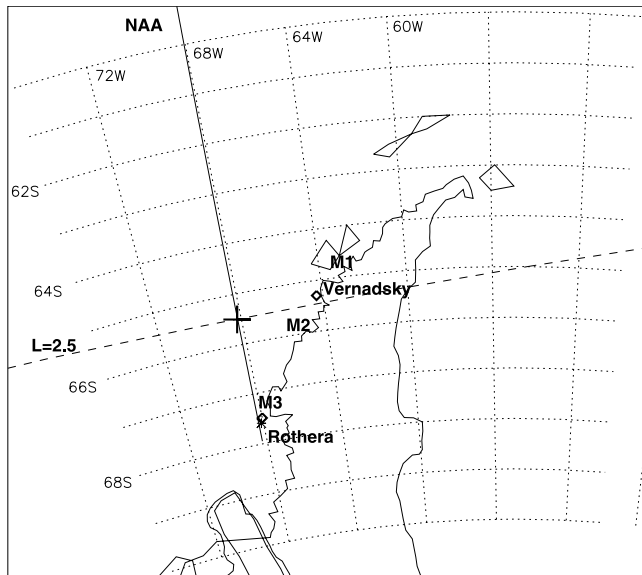
<sup>9</sup>Space Instrumentation Group, Department of ACSE, University of Sheffield, Sheffield, UK.

<sup>10</sup>Department of Physics, University of York, York, UK.

<sup>11</sup>STAR Laboratory/Electrical Engineering Department, Stanford University, Stanford, California, USA.

### 1. Introduction

[2] It has previously been shown that the plasma mass density in the inner magnetosphere can be determined from the interstation phase and amplitude spectra of geomagnetic pulsations between pairs of ground-based magnetometers [Menk *et al.*, 1999, 2000]. It is also well known that the plasmaspheric electron number density can be determined by analysing the propagation of ducted man-made whistler-mode signals [Clilverd *et al.*, 2000]. Both of these measurements tend to be single point, or at least single longitude. The pulsations are essentially a daytime probe, while the whistler-mode signals are more often observed during nighttime. Good spatial information on the structure of the inner magnetosphere can be obtained using satellite-based experiments, although the measurement repetition rate tends to be low.



**Figure 1.** Map of the Antarctic Peninsula showing the geographic positions of the low-power magnetometer sites (M1, M2, M3), Vernadsky magnetometer (diamond), the VLF Doppler receiver site (asterisk), and the typical field line footprint of study (cross).

[3] The recently launched IMAGE satellite is able to image the structure of the entire plasmasphere [Sandel *et al.*, 2000, 2001]. The images are necessarily line-of-sight and require a degree of interpretation to deconvolve the multilayered information. Crucially, information on the ion composition of the inner magnetosphere is necessary to accurately interpret the images during varying periods of geomagnetic activity levels.

[4] Dent *et al.* [2003] compared magnetometer/VLF/Image RPI measurements made at different latitudes and longitudes for one specific time. In this study we provide the first analysis of data that describes the structure of the inner magnetosphere using all of these techniques, i.e., pulsations, ducted whistler-mode signals, and in situ satellite measurements, observing the same point on a geomagnetic field-line and at the same longitude. We discuss two representative days and show that the plasma number density results determined from the different techniques are consistent, and can be used to intercalibrate the EUV Imager images produced by the IMAGE satellite. A coupled plasmasphere-ionosphere model is compared against the observations as a further intercalibration.

## 2. Experimental Details

### 2.1. Magnetometer Observations

[5] In order to monitor the plasma mass density at  $L = 2.5$ , we examined ULF geomagnetic field line resonances (FLRs) during January and February 2001 using a combination of ground magnetometer arrays. The primary magnetic data were obtained with an array of 3 Low-Powered Magnetometers (LPMs) deployed and operated by the British Antarctic Survey since December 2000 specifically for this project. Figure 1 shows the locations of these closely spaced stations located on the Antarctic Peninsula

and spanning  $2.3 < L < 2.8$  ( $63^\circ\text{S} - 68^\circ\text{S}$  and  $63^\circ - 68^\circ\text{W}$  geographic). Their relative locations are indicated by labels M1, M2, and M3.

[6] Each LPM is an autonomous three-component fluxgate magnetometer system developed for remote operation in temperatures down to  $-70^\circ\text{C}$ , in an environment of extreme gales and high static discharge hazard. The systems use GPS timing, solar panels, and rechargeable batteries, and they conserve power by switching the sensors off between samples. Each system is visited just once per year to download data. The LPM data were sampled in geographic coordinates at 0.1 Hz (averaged from the 5 s measurements) with  $<1$  nT resolution, and rotated into geomagnetic (H, D) coordinates before analysis.

[7] An additional magnetometer at the Ukrainian Antarctic station Vernadsky ( $65^\circ\text{S}$ ,  $64^\circ\text{W}$ ; see Figure 1) provided data from close to the midpoint of two of the LPM systems. This comprises a three-component LEMI-008 type fluxgate magnetometer that samples the geomagnetic field each 1 s with 0.1 nT resolution. This magnetometer has transmitted its data to the British Geological Survey INTERMAGNET project since February 2001.

[8] Additional information on the variation of resonant frequency (and hence mass density) with latitude was obtained using data from the combined SAMNET and BGS magnetometer arrays obtained from the University of York, UK. For this purpose 5 s measurements with  $<0.1$  nT resolution were examined from 7 SAMNET/BGS magnetometer stations spanning  $L = 2.24$  to  $L = 5.47$  between  $0^\circ\text{W}$  and  $21^\circ\text{W}$  geographic longitude.

[9] Techniques for the detection of FLRs were summarized by Menk *et al.* [1999, 2000]. When examining data from latitudinally separated magnetometers, the resonant frequency is identified by the peak in H-component cross-power and cross-phase, and a unity crossing in H-component power ratio, approximately midway between the stations. Where only one station is available the resonance is indicated by a peak in the power ratio H/D and a rapid change in polarization, i.e., in the phase between the H and D components.

[10] The diurnal pattern of resonance behavior is most clearly seen with dynamic cross-phase spectra [e.g., Waters *et al.*, 1991, 1995]. However, all measurements of resonant frequency in this paper were obtained from inspection of discrete FFT cross-power, coherence, power difference, power ratio, and cross-phase spectra on selected days over time windows between 30 and 60 min long, the shorter windows being used under more active conditions.

[11] Mass densities were calculated from the measured resonant frequencies using the analytical expressions described by Taylor and Walker [1984] and Walker *et al.* [1992]. These assume decoupled toroidal mode oscillations and yield essentially identical results to the models described by Orr and Matthew [1971]. We used a model distribution of the form  $R^{-3}$ , where  $R$  is the radial distance in the equatorial plane, since this is where the magnetic field strength, and hence Alfvén velocity along a field line generally reach a minimum. This is invalid for very low or high latitude field lines because of mass loading by ionospheric heavy ions and field line distortion, but it is a satisfactory approximation at  $L = 2.5$ . The choice of power law describing the radial variation in density within the

plasmasphere has been discussed by, e.g., *Menk et al.* [1999] and *Denton et al.* [2001], although the calculated mass density is relatively insensitive to the actual power law [Orr and Matthew, 1971; Walker et al., 1992].

[12] The uncertainty in our calculated mass densities (20–30%) depends mainly on uncertainty in the frequency measurement with uncertainties typically of order 10–15%. *Menk et al.* [1999] discussed the relationship between these two uncertainties and found the mass density uncertainty to be typically double the uncertainty in frequency measurement. We have assumed a dipole magnetic field, and at  $L = 2.5$  this introduces negligible error. The equatorial mass density can also be determined from comparison of resonance harmonics [Schulz, 1996]. That method requires observed frequencies to be accurate to 6% in order to obtain equatorial mass density with a precision of 30% [Denton and Gallagher, 2000]. The guided poloidal mode eigenfrequency is  $\sim 30\%$  lower than our assumed purely toroidal mode oscillations so that wave mode coupling may cause the mass density to be slightly underestimated.

## 2.2. VLF Data

[13] Equatorial electron number densities ( $N_{eq}$ ) were estimated for  $L = 2.5$  using man-made VLF whistler-mode signals recorded at Rothera in Antarctica (geographic 67.5°S, 68.1°W,  $L = 2.8$ ). The VLF Doppler experiment at Rothera receives ducted whistler-mode signals from the east coast U.S. Navy transmitter NAA, 24.0 kHz, 1 MW at Cutler, Maine (geographic 44.6°N, 67.3°W,  $L = 3.0$ ). Narrowband receivers of the type described by Thomson [1981] are able to separate the whistler mode signals from the stronger subionospheric signal and measure the group delays ( $t_g$ ), Doppler shifts, and arrival bearings of the whistler mode component. Results from these experiments have been published by Saxton and Smith [1989] and Clilverd et al. [2000].

[14] The method of determining the group delays,  $L$  shells, and propagation paths of the whistler mode signals was discussed in detail by Smith and Clilverd [1991] and Clilverd et al. [1991]. The group delay times of the whistler mode signals are determined by cross correlating the plasmaspheric signal with the subionospheric signal, accumulating the coefficients for 15 min at a time. The  $L$  shell has in the past been determined by measuring the difference in group delay time for two different transmitter frequencies traveling in the same duct. The time difference is due to dispersion along the field line and is dependent on the  $L$  shell of propagation. In this study we use the simplification that the average group delay observed is representative of equatorial electron number density conditions at  $L = 2.5$ . The error incurred by using this simplification is  $\sim 0.1 L$  [Saxton and Smith, 1989]. As the longitude of the transmitter and the receiver are so similar the whistler-mode signals exit the ionosphere close to the Antarctic Peninsula LPM array. The intercomparison between this technique and FLRs has been discussed by *Menk et al.* [1999] and *Dent et al.* [2003].

## 2.3. IMAGE Satellite Data

[15] The IMAGE spacecraft is in an elliptical polar orbit with an apogee altitude of 7.2 Earth radii (46,000 km) and a perigee altitude of 1000 km, and completes one orbit

every 14.2 hours. This study presents data from two of the seven instruments on board the spacecraft. The RPI experiment provides the electron number density as IMAGE transits through the plasmasphere, and the EUV Imager experiment images geospace using  $\text{He}^+$  emissions when the spacecraft is near apogee looking back toward the Earth.

[16] The RPI instrument [Reinisch et al., 2000; Carpenter et al., 2002] is a low-power radar which operates in the radio frequency bands which contain the plasma resonance frequencies characteristic of the Earth's magnetosphere (3 kHz to 3 MHz). In the passive receive-only mode, RPI receives the natural plasma wave emissions from 3 kHz to 1.1 MHz, using these to determine the electron number density to within an accuracy of 10%. In this study we are able to make use of the fact that the orbit of IMAGE takes it close to the geomagnetic equator at  $\sim L = 2.5$  so that there is no adjustment to make in order to compare the RPI measurements with the VLF Doppler electron number density estimates.

[17] The Extreme Ultraviolet Imager (EUV Imager) images the distribution of  $\text{He}^+$  in Earth's plasmasphere by detecting its resonantly scattered emission at 30.4 nm [Sandel et al., 2001]. Effective imaging of the plasmaspheric  $\text{He}^+$  requires global snapshots in which the high apogee and the wide field of view of the EUV Imager provide in a single exposure a map of the entire plasmasphere. The 30.4 nm feature is relatively easy to measure because it is the brightest ion emission from the plasmasphere, it is spectrally isolated, and the background at that wavelength is negligible. Line-of-sight measurements are easy to interpret because the plasmaspheric  $\text{He}^+$  emission is optically thin, so its brightness is directly proportional to the  $\text{He}^+$  column abundance. The EUV Imager instrument consists of three identical sensor heads, each having a field of view of 30°. These sensors are tilted relative to one another to cover a fan-shaped field of  $84^\circ \times 30^\circ$ , which is swept across the plasmasphere by the spin of the satellite. EUV Imager's spatial resolution is  $\sim 0.6^\circ$  or  $\sim 0.1 R_E$  in the equatorial plane seen from apogee. The sensitivity is sufficient to map the position of the plasmopause with a time resolution of 10 min or better.

## 2.4. Plasmasphere-Ionosphere Model

[18] We compare the measured densities with modeled values of electron and ion number densities in the equatorial region of the  $L = 2.5$  field-line at Antarctic Peninsula longitudes. These have been determined from SUPIM, the Sheffield University Plasmasphere Ionosphere Model [Bailey and Selleck, 1990; Bailey et al., 1993; Bailey and Balan, 1995]. In SUPIM, coupled time-dependent equations of continuity, momentum, and energy balance are solved along closed magnetic field lines between altitudes of around 120 km in conjugate hemispheres for the concentrations, field-aligned fluxes, and temperatures of the  $\text{O}^+$ ,  $\text{H}^+$ ,  $\text{He}^+$ ,  $\text{N}_2^+$ ,  $\text{O}_2^+$ , and  $\text{NO}^+$  ions and the electrons. For the present study the geomagnetic field is represented by an eccentric dipole. The concentrations and temperatures of the neutral gases are obtained from the MSIS86 thermospheric model [Hedin, 1987], the neutral wind velocities from the HWM90 neutral wind model [Hedin et al., 1991], and the solar EUV fluxes from the EUVAC solar EUV flux model of Richards

*et al.* [1994]. The remaining model inputs, such as the photoionization and photoabsorption cross sections, the chemical reaction rates, the heating and cooling rates, and the collision frequencies are described by *Bailey and Selleck* [1990], *Bailey et al.* [1993], and *Bailey and Balan* [1995].

### 3. Results

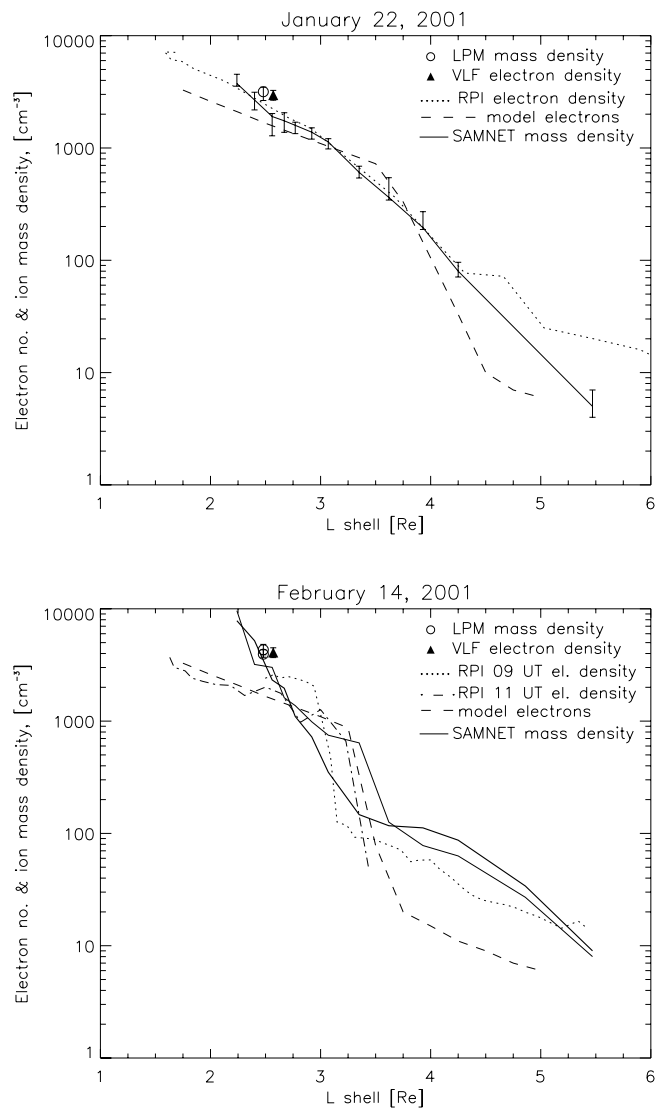
#### 3.1. Overview

[19] The low-power magnetometers were deployed in Antarctica and commenced data recording in December 2000 and have operated continuously since then. The VLF Doppler experiment has been operating concurrently. This paper presents results from 2 days of data from the first summer of operation, through to early March 2001. Geomagnetic conditions during this time were relatively quiet, with no significant magnetic storms. The two days selected for detailed study, 22 January and 14 February 2001, are two of the more disturbed periods during this time, when the occurrence of ducted VLF signals and ULF pulsation activity is likely to peak. To assess this, we examined (1) monthly scatter plots of the group delay of NAAVLF Doppler signals received at Rothera, (2) monthly scatter plots of the local time distribution of these Doppler signals, and (3) time series records from the colocated magnetometers.

[20] On the first day examined, 22 January,  $K_p$  reached 4- after peaking the previous day at 4<sub>0</sub>.  $A_p$  reached 19 on 21 January, again indicative of moderately disturbed conditions, while the previous 10 days had registered a maximum of only 9. The maximum  $K_p$  on the second selected day, 14 February, was again 4-, after peaking at 5- just before midnight on 13 February.  $A_p$  reached 25 on 13 February, again indicative of moderately disturbed conditions, while the previous 10 days had registered a maximum of 14 on 6 February. The sunspot number on 22 January was 93 while for 14 February it was 68. Thus both selected days represent moderately disturbed conditions, occurring after an extended quiet period.

[21] The RPI electron number density data can also provide information on the state of the plasmasphere during these days. Figure 2 shows the scaled plasma wave results for passes through the plasmasphere made at about 0800–1100 UT on both study days (dotted and dash-dot lines), pulsation FLR results from the U.K. SAMNET array (solid lines), and the empirical model results of *Carpenter and Anderson* [1992] (dashed lines). The RPI results include extrapolations from the point of measurement to the geomagnetic equatorial plane at each  $L$ -shell. This is done by simply assuming the density along a field line is uniform over the small range of  $L$ -shells considered here where the satellite is close to the geomagnetic equator [*Goldstein et al.*, 2001]. Also included in the figures are mass densities and electron densities derived from the LPMs and VLF experiments near  $L = 2.5$  on the Antarctic Peninsula (open circles and filled triangles, respectively). For the 22 January plot, uncertainty in the ULF-derived mass density is indicated by error bars. For clarity these have not been shown in the 14 February plot, although the uncertainties are similar.

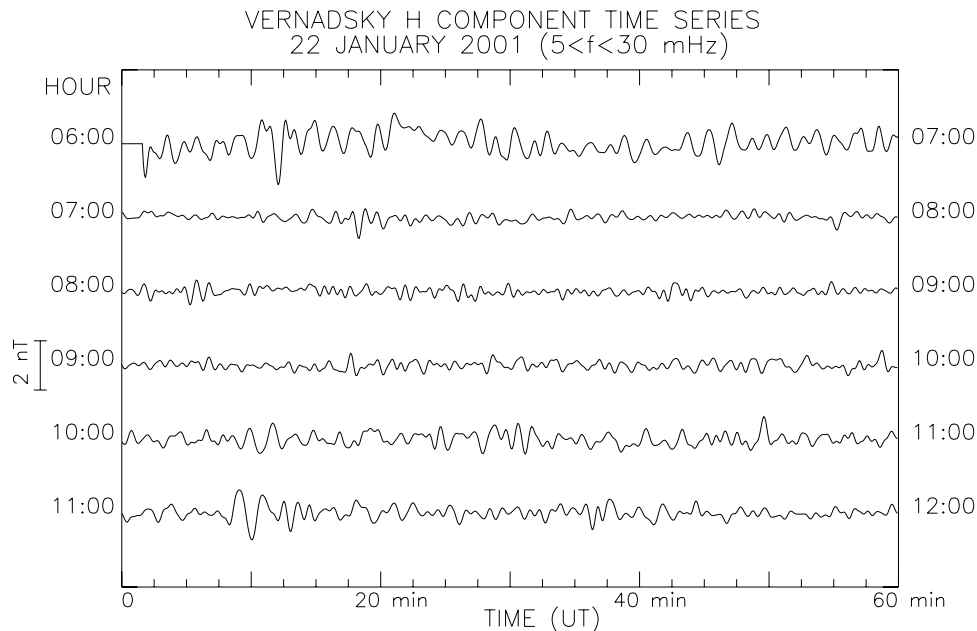
[22] The 22 January data (top panel) are best represented by an extended plasmopause about 0.5  $L$  wide situated somewhere near  $L = 3.8$ . The IMAGE spacecraft pass is



**Figure 2.** Variation of plasma density through the magnetosphere near 1000 UT on 22 January (top panel) and 14 February 2001 (bottom panel). IMAGE RPI electron densities are indicated by the dotted line (0800 UT on 22 January and 0900 UT on 14 February) and dash-dot line (1100 UT on 14 February). Ion mass densities determined using the SAMNET/BGS magnetometer arrays are denoted by the solid line (0800–0850 UT on 22 January and 0600–0650 UT on 14 February) and heavy solid line (0800–0850 UT on 14 February). The dashed line in both plots represents the empirical model of *Carpenter and Anderson* [1992] for these geomagnetic conditions. Finally, open circles and solid triangles indicate  $L = 2.5$  ion mass density and electron density at the Antarctic Peninsula.

from 0802 to 0845 UT and 0438 to 0502 MLT. The model and RPI electron densities are similar around  $L = 2.5$ – $3.3$  even though there is little overall agreement throughout the plasmasphere. Beyond  $L = 4.3$  the RPI density extrapolations become unreliable as the spacecraft ascends to  $>20^\circ$  latitude above the equatorial plane.

[23] The plasma mass density (solid line) was determined between 0800 and 0840 UT for a range of  $L$ -values using SAMNET and BGS magnetometer data at about  $0^\circ$



**Figure 3.** An example of the Vernadsky magnetometer data on 22 January, 0600–1200 UT, showing small pulsations of about 1 nT amplitude.

geographic longitude. There is good overall agreement between the RPI experiment at  $70^\circ\text{W}$  geographic longitude and the mass density values, although the mass densities tend to be slightly lower. This is primarily because of the influence, at this time of year, of the longitudinal amplitude of the annual variation in plasmaspheric electron number densities [Clilverd *et al.*, 1991]. At longitudes of the Antarctic Peninsula ( $60$ – $70^\circ\text{W}$  geographic), high concentrations of ion and electrons would be expected in January compared with those in June, while at the longitudes of Europe, Africa, and India such high concentrations would not be expected in January.

[24] On 14 February (bottom panel) a reasonably well-defined plasmopause is present in the IMAGE RPI measurements. Data from two orbits are shown in the figure: 0901–1018 UT (1516–1504 MLT; dotted line), and 1101–1157 UT (0349–0293 MLT; dot-dash line). The plasmopause location is around  $L = 2.97$ – $3.12$  and  $L = 3.22$ – $3.5$  during these passes. Ion mass density data determined from the SAMNET/BGS magnetometer array are also shown for two times: 0600–0650 UT (solid line) and 0800–0850 UT (heavy line). There are significant differences in the shape and position of the plasmopause derived from the measurements at these times. This suggests that there are temporal and longitudinal variations in both electron density and heavy ion concentration during this day.

[25] Again, the  $L = 2.5$  ion mass density and electron density measurements are somewhat higher than values at other longitudes because of the longitudinal amplitude of the annual variation in plasmaspheric number density. Note that at  $L < 2.7$  the ion mass densities are also significantly higher than the RPI and model electron densities. In the case of the model values this is because of the lack of any annual variation effects.

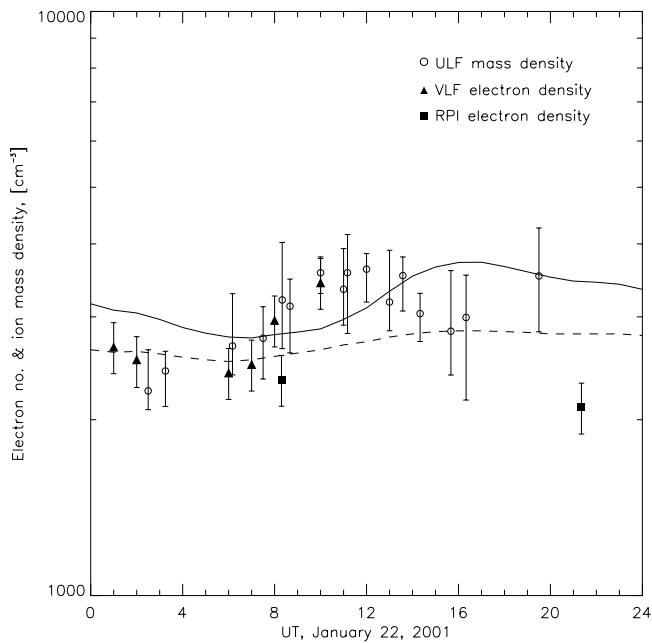
[26] Clearly, when comparing the results from a range of instruments it is important to consider their relative loca-

tions. This paper uses data from longitudes in the Antarctic Peninsula region. On 22 January 2001 IMAGE passes the geomagnetic equator near  $L = 2.5$  between  $290$  and  $310^\circ\text{E}$  geographic. Thus the ground-based and satellite measurements are closely collocated. On 14 February 2001 IMAGE passes the geomagnetic equator near  $L = 2.5$  at about  $240^\circ\text{E}$  (geographic) or  $50^\circ$  west of the Antarctic Peninsula region where the ground-based measurements are made. Although still in the same longitudinal region in terms of the likely influence of the annual variation in plasma density at  $L = 2.5$  the difference may still be significant as further analysis shows.

### 3.2. 22 January 2001

[27] The ground-based Antarctic Peninsula magnetometer data come from a region centered on the Ukrainian Station of Vernadsky. The determination of the magnetospheric plasma mass density distribution is dependent on the presence of field-line resonances in the ULF data. Representative magnetometer time series from 0600–1200 UT from Vernadsky are shown in Figure 3, where six 1-hour sections of data are stacked vertically. The amplitude scale is 2 nT and is shown on the left hand y-axis. The plot shows that small but discernable pulsations occur at most times, allowing good UT coverage for the mass density analysis. This is expected for moderately disturbed geomagnetic conditions and was a primary influence on the dates selected in this study.

[28] The diurnal variation in magnetospheric plasma mass density near the equatorial plane at  $L = 2.5$  on 22 January 2001 is shown by the open circles with error bars in Figure 4. The average mass density is close to  $3000$   $\text{amu cm}^{-3}$  for most of the day, although slightly lower values occur during the nighttime hours before 0800 UT (0400 LT). Reasonable data coverage extends until around 1700 UT when pulsation activity became weaker. Also



**Figure 4.** Ion mass density (open circles) and electron density (triangles) variations on 22 January 2001. SUPIM model results for electron number density (dashed line) and ion mass density (solid line) are shown for comparison. RPI electron number density values are given by the solid squares. Error estimates for each observation are indicated by the vertical lines and are typically 10%.

shown on Figure 4 are the electron number densities, derived by both the VLF Doppler experiment (filled triangles with error bars) and the RPI experiment (filled squares with error bars). These show a reasonable correlation with the mass densities during the first half of the day, when VLF data were available.

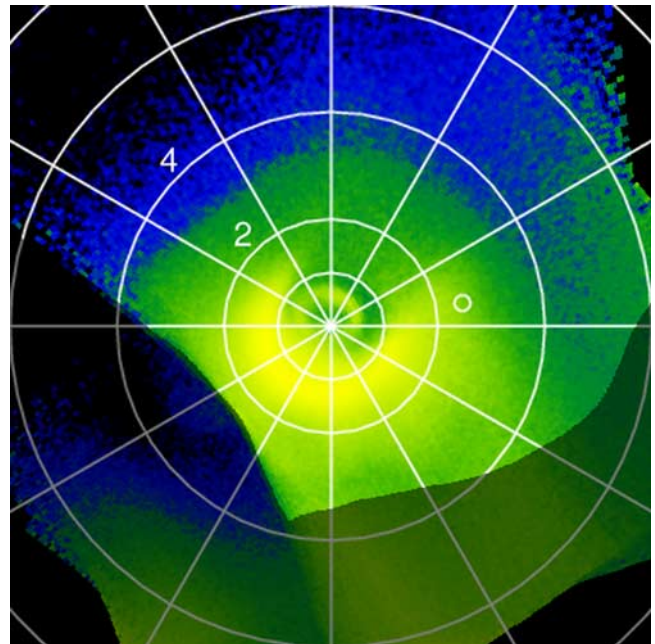
[29] Figure 4 also shows results from the SUPIM model at the geomagnetic equator of the  $L = 2.5$  field-line based on the solar and geomagnetic conditions for this day. The model starts from arbitrary initial conditions and then runs for several calculation days until the output values converge. The dashed line indicates the diurnal variation of the model electron number density, which shows good agreement with the equivalent VLF values during the nighttime (0000–0700 UT). The solid line indicates the plasma mass density, which also shows reasonable agreement with the ULF observations during the first third of the day, and a similar diurnal range. However, the phase of the model variation is offset from that of the observations by 4–5 hours. The difference between the model electron density and model mass density results is due to the presence of heavy ion populations in the model. During the day the proportion of  $\text{He}^+$  to  $\text{H}^+$  in the model is about 2.6% by number and 10.5% by mass, while  $\text{O}^+$  increases from 0.7% to 1.0% by number or 11–16% by mass.

[30] Figure 5 is an EUV Imager image from 22 January 2001. It is the sum of eight individual EUV Imager images that have been rebinned in  $L$ -shell and magnetic longitude space [Sandel *et al.*, 2003] and then added. Thus the brightness of this image represents the  $\text{He}^+$  distribution in

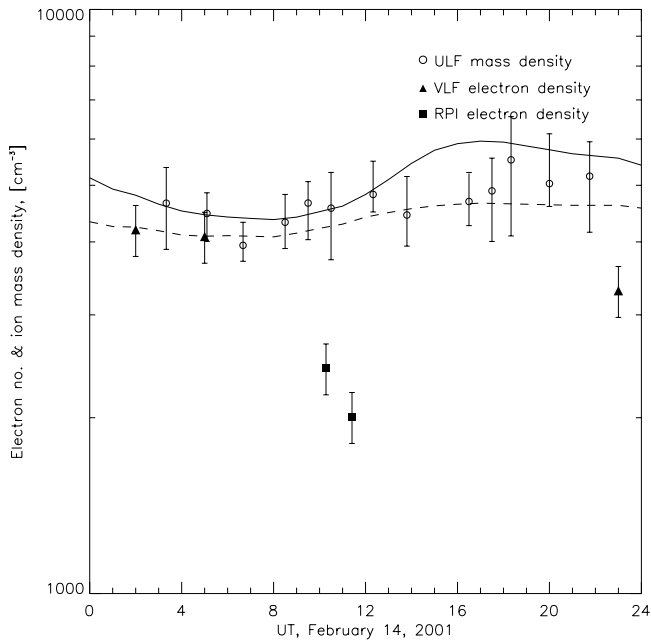
the plane of the magnetic equator. The innermost white circle corresponds to the solid Earth. Auroral emission is present about half-way from the center to the Earth's limb between magnetic longitudes of approximately 0 and 180°. Earth's shadow moves in magnetic longitude, so it is smeared by summing the rebinned images. The position of the field-line at  $L = 2.5$ , magnetic longitude = 10°E (shown by the white circle to the right of the Earth) is outside the shadow for all the images used in this work, so our determination of  $\text{He}^+$  column abundances is not compromised. We used this type of plot to determine the line-of-sight  $\text{He}^+$  abundance for the line of sight at the white circle and found a value of  $5.8 \times 10^{10} \text{ cm}^{-2}$ .

### 3.3. 14 February 2001

[31] The magnetospheric plasma mass density distribution near  $L = 2.5$  on 14 February 2001 is represented by the open circles in Figure 6. The plot format is the same as Figure 4 for 22 January. It is noticeable, however, that on 14 February the mass densities are somewhat higher, at about 4500 a.m.u.  $\text{cm}^{-3}$ , but the electron number densities determined from the VLF Doppler experiment still track the mass densities. However, the RPI ‘snapshot’ values (2200 el.  $\text{cm}^{-3}$ ) are significantly lower, by a factor of 2 around 1000 UT. As mentioned previously, near 0900 UT and 1100 UT the IMAGE satellite was respectively about 150° east and 50° west of the ground stations, and this may



**Figure 5.** EUV Imager view of the near-Earth environment on 22 January 2001. This composite includes eight individual EUV images acquired from 0051 to 0202 UT, when IMAGE was at magnetic latitudes from 77°N to 68°N and radial distances from 8.11 to 7.90 Re. The image is projected into the plane of the magnetic equator, and the white grid shows  $L$  and magnetic longitude, with 0° magnetic longitude to the right. The white circle is at the position of the Antarctic Peninsula  $L = 2.5$  field-line. The dark mask over the lower parts of the image marks regions of incomplete coverage or scattered sunlight.



**Figure 6.** Ion mass density and electron density on 14 February 2001, in the same format as Figure 4.

be why the in situ observations are significantly different from the other experiments. Analysis of naturally occurring whistlers propagating at about  $320^{\circ}\text{E}$  geographic longitude confirms electron number densities of  $4000 \text{ el. cm}^{-3}$  at  $L = 2.5$ , while estimates from a second VLF Doppler system in Dunedin, New Zealand [Thomson, 1981] also agree with the high electron values observed on this day.

[32] Figure 6 also shows results from the SUPIM model at the geomagnetic equator of the  $L = 2.5$  field line based on the solar and geomagnetic conditions for this day. The dashed line again indicates the diurnal variation of the model electron number density, which shows good agreement with the equivalent VLF values prior to 0800 UT. The solid line indicates the plasma mass density, which also shows good agreement with the ULF observations, especially during the first half of the day. A significant difference between this model run and the one undertaken for 22 January 2001 above is that the photoionization rate used for 22 January was artificially made half as large in order to produce similar density results to those observed. The photoionization rate used on 14 February was typical for standard model runs, while those for 22 January represent lower than normal rates. It is unclear how the difference in photoionization rate between the two days is driven physically. The percentage of  $\text{He}^+$  to  $\text{H}^+$  in the model for 14 February is again about 2.6% by number and 10.5% by mass, while  $\text{O}^+$  increases throughout the day from 0.7% to 1.0% by number or 11–16% by mass.

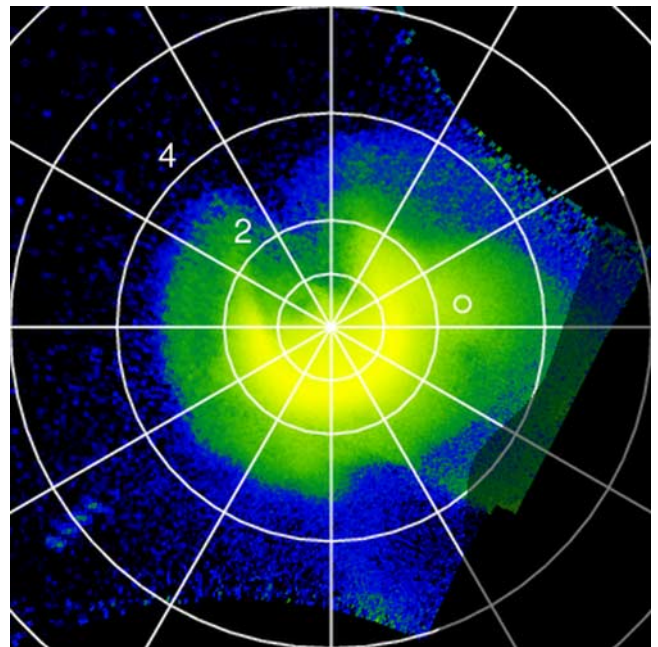
[33] The EUV Imager data for 14 February were analyzed when IMAGE was close to apogee on two orbits. The first measurements at  $L = 2.5$  were made between 2344 (the previous day) and 0126 UT and gave an average column abundance of  $6.7 \times 10^{10} \text{ cm}^{-2}$ . The second period was during 2142–2213 UT, where the column abundance was  $11.0 \times 10^{10} \text{ cm}^{-2}$ . This second period appears to be

associated with a tongue of more intense  $\text{He}^+$  emission as shown in Figure 7, where the location of the geomagnetic equator at  $L = 2.5$  at the longitude of study in this paper is indicated by the white circle. Compared with the model values the ULF mass densities do not appear to show any evidence of such a significant increase during the period around 2100 UT. However, compared with the VLF electron number densities observed at 2300 UT significant levels of heavy ions could be likely.

#### 4. Discussion

[34] In previous studies of the ion and electron structure of the plasmasphere, Menk *et al.* [1999] and F. W. Menk *et al.* (Monitoring the plasmopause using geomagnetic pulsations: 1. Field line resonances, submitted to *Journal of Geophysical Research*, 2003). used spatially separated measurements, correcting for local time effects in order to intercompare the results. In this study we show the advantage of using colocated instruments in that no corrections need to be made in order to compare the results.

[35] Here we have chosen 2 days with similar geomagnetic conditions, where the IMAGE satellite passed close to the equatorial region of an  $L = 2.5$  field-line connecting to a cluster of ground-based experiments measuring similar parameters. While the 2 study days seem to relate to similar conditions within the inner plasmasphere, the outer plasmasphere (taken to be  $L > 3.0$  here) is quite different on the 2 days, with more structure and a well defined plasmopause present on 14 February 2001. This may be due to different



**Figure 7.** EUV Imager view of the near-Earth environment at 2145 UT on 14 February 2001. This composite includes four individual EUV images acquired from 2142 to 2213 UT, when IMAGE was at magnetic latitudes from  $61^{\circ}\text{N}$  to  $56^{\circ}\text{N}$  and radial distances from 6.44 to 5.92 Re. The presentation is the same as in Figure 5.

solar wind conditions. The solar wind speed was almost constant around  $320 \text{ km s}^{-1}$  over most of 20–21 January 2001, reaching a plateau around  $500 \text{ km s}^{-1}$  on 22 January. The solar wind speed was in the range  $400\text{--}450 \text{ km s}^{-1}$  on 12 February but was in the range  $480\text{--}610 \text{ km s}^{-1}$  on 13 February and 14. The solar wind proton density was between  $2$  and  $20 \text{ cm}^{-3}$  over 20–22 January, with a short large increase to  $60 \text{ cm}^{-3}$  late on 21 January. Over 12–13 February the solar wind proton density was generally between  $5$  and  $10 \text{ cm}^{-3}$  and  $<6 \text{ cm}^{-3}$  on 14 February. Thus solar wind pressure (which is proportional to  $nV^2$ ) was significantly higher over 12–13 February than over 20–21 January. This may also explain why double the photoionization rates were required in the SUPIM model to reproduce the ground-based observations on 14 February compared with 20 January. The IMF  $B_z$  component was similar during both 20–21 January and 12–13 February, with positive variations and several excursions to below  $-7 \text{ nT}$ .

[36] Using the EUV Imager experiment on IMAGE, we find that the column abundance of  $\text{He}^+$  changed from  $5.8$  to  $6.7 \times 10^{10} \text{ cm}^{-2}$  between 22 January and 14 February. This is in qualitative agreement with the observed change in general plasma concentration between the 2 days. More generally, it appears that a typical value of  $6 \times 10^{10} \text{ cm}^{-2}$  is appropriate for comparison with the ground-based observations. We used this value to investigate the unexpectedly lower electron number density values on 14 February when IMAGE was  $50^\circ$  west of the ground stations. EUV Imager values from that region made at 1300–1400 UT indicate column abundances of only  $3.6 \times 10^{10} \text{ cm}^{-2}$ , which is suggestive of a region of plasma depletion that may have been present during this period and that could explain the IMAGE RPI values determined only 3 hours earlier.

[37] The ratio of  $\text{He}^+$  to  $\text{H}^+$  ions in the plasmasphere can be described from an empirical expression given by Craven *et al.* [1997]. The ratio decreases with radial distance and depends on the solar activity  $10.7 \text{ cm flux}$ . At  $L = 2.5$  the typical mass ratio of  $\text{He}^+:\text{H}^+$  is about 0.1, with a range from 0.03 to 0.3. The solar flux measure F10.7 was typically about 150 during January/February, which Craven *et al.* [1997] suggests should result in the abundance of  $\text{He}^+:\text{H}^+$  being 10% at  $L = 2.5$ . Using the values of electron number density and ion mass density from 0800 UT on 22 January 2001 and assuming charge neutrality, we find that  $\text{He}^+$  is about 3.8% of the abundance of  $\text{H}^+$ . Similar constraints on the values determined from 0400 UT on 14 February 2001, when the electron number density is generally higher on this day than on 22 January ( $4000 \text{ el. cm}^{-3}$  compared with  $3000 \text{ el. cm}^{-3}$ ), gives  $\text{He}^+$  abundance levels of about 3.3% by number and about 13–14% by mass. However, the estimated number density of  $\text{He}^+$  atoms actually increases from 110 to  $130 \text{ cm}^{-3}$ , or a ratio of about 1.2. From the EUV Imager images we know that the abundance of  $\text{He}^+$  on the first half of these two days also shows a similar increase ( $6.7/5.8 = 1.2$ ).

[38] Having determined the typical column abundance figure for  $L = 2.5$  can we relate this to the concentration values that we find in the equatorial region? Comparison between space-based and ground-based systems are reasonable when the plasma is distributed symmetrically and uniformly about the geomagnetic equator. Tu *et al.* [2003]

showed that this was the case. Modeling also shows that the measured radial profiles of  $\text{He}^+$  column density near  $L = 2.5$  are similar in shape to those computed by integrating the RPI electron density profiles such as those shown in Figure 2 [Sandel *et al.*, 2003]. This implies that we can take the  $L$ -dependence for  $\text{He}^+$  near  $L = 2.5$  to be the same as that measured for electrons by RPI. We can then use the profiles of electron density measured by RPI to compute an effective path length ( $P$ ). Here  $P$  is the length of the column of  $\text{He}^+$  having uniform density equal to the region of maximum, i.e., equatorial, density along the line of sight, and having the measured column density. For the plasma distributions seen near  $L = 2.5$ , we find  $P = 1.8 \text{ Re}$ . Then the measured column density corresponds to an equatorial concentration of  $52 \text{ He}^+ \text{ cm}^{-3}$  in this region. With  $\text{H}^+$  levels close to the electron number density of about  $3000 \text{ el. cm}^{-3}$  we find that this would result in  $\text{He}^+$  occurrence levels of 1.7% by number and 6.9% by mass. As the VLF and ULF measurements are biased to about 10% of the length of the magnetic field line that lies near the equatorial plane ( $0.5 R_E$  for an  $L = 2.5$  field-line) we could probably multiply the above intercomparison result by a factor of 2 (a value suggested by Sandel *et al.* [2003]) to allow for the difference in effective path length, i.e.,  $1.8 \text{ Re}$  to  $0.5 \text{ Re}$ . As a result we might expect about 3–4%  $\text{He}^+$  by number and 10–15% by mass from these calculations. Although further modeling will be required to complete this picture in detail it does appear that this simple calculation confirms the analysis made from the ground-based observations.

[39] At 2200 UT on 14 February the IMAGE EUV Imager measurements suggest that the abundance of  $\text{He}^+$  exceeded the values that were observed on 22 January by almost a factor of 2, i.e., 11 instead of  $6 \times 10^{10} \text{ cm}^{-2}$ . Although data on the electron number densities are sparse at this time, the disparity between the ULF mass densities and the VLF electron number densities between 2000 and 2300 UT shown in Figure 6 suggests that heavy ions must be an influence. Once again requiring charge neutrality we find that the 2000–2300 UT observations are consistent with  $\text{He}^+$  being about 20% of  $\text{H}^+$  by number, with the actual levels being about  $540 \text{ cm}^{-3}$ . As a ratio of the values earlier in the day this is  $540/130 = 4.9$ , which is significantly higher than the ratio of the equivalent column abundances, i.e.,  $11/6.7 = 1.6$ . A near doubling of just  $\text{He}^+$  in the tongue of emission visible at 2100 UT would not be enough to produce the majority of the change in mass density levels to above  $5000 \text{ a.m.u cm}^{-3}$  observed at about 2100 UT. This suggests the presence of heavier ions such as  $\text{O}^+$ . Further examples will need to be studied in order to clarify the relationship between ions and electrons in this type of feature.

[40] With some confidence in intercomparing the EUV Imager and ground-based experiments we can now consider the relevance of the EUV Imager observations made at the location where the RPI measurements were significantly lower than the ground station observations on 14 February. EUV Imager levels of  $3.6 \times 10^{10} \text{ cm}^{-2}$  are a factor of about  $3.6/6.7 = 0.53$  down on the Antarctic Peninsula values. As we estimated that electron number densities there were  $4000 \text{ el. cm}^{-3}$  we would estimate RPI values of  $2150 \text{ el. cm}^{-3}$  to the west. Clearly, this value equates well to the observed RPI value of  $2200 \text{ el. cm}^{-3}$  made at about



1100 UT and confirms the presence of significant depletion structures in the plasmasphere on this day.

## 5. Conclusions

[41] Using two ground-based experiments and two satellite-borne experiments, we have been able to interpret changes in plasmaspheric composition at the same point in space during moderate geomagnetic activity. We have analysed data from 22 January 2001 and 14 February 2001 on an  $L = 2.5$  magnetic field line at the longitude of the Antarctic Peninsula. Mass density was determined from an array of magnetometers, while the electron number density was determined from analysis of the group delay of man-made VLF transmissions from north east America. The IMAGE satellite RPI experiment provided in-situ plasma wave measurements of the electron number density in the equatorial region of the  $L = 2.5$  field line, while a few hours later the IMAGE EUV Imager experiment was able to resolve the  $\text{He}^+$  abundance by looking back toward the same spot.

[42] On 22 January 2001 the mass density and electron number density were very similar. The levels were consistent with a moderately disturbed plasmasphere. The electron number density from IMAGE RPI was also in good agreement with the ground-based measurements at the times when it was close to the equatorial region of the  $L = 2.5$  field line. The IMAGE EUV Imager measurements of  $\text{He}^+$  were relatively uniform throughout the plasmasphere, again indicative of moderately disturbed conditions, and the column abundance value of  $5.8 \times 10^{10} \text{ cm}^{-2}$  is taken to represent those conditions. There was about 4% by number and 16% by mass of  $\text{He}^+$  relative to  $\text{H}^+$  at this time. On 14 February 2001 there appeared to be significantly more response of the plasmasphere to the moderate ( $K_p = 5$ ) activity levels. The electron number density and the mass density showed a significant increase compared with January 22 for most of the day. The IMAGE EUV Imager experiment also indicated that a 20% increase in the abundance of  $\text{He}^+$  could be seen between these two days.

[43] We find that for the longitudes of the Antarctic Peninsula and therefore the Americas, an EUV Imager  $\text{He}^+$  column abundance value of  $6 \times 10^{10} \text{ cm}^{-2}$  is equivalent to plasmaspheric electron density levels of  $3000 \text{ cm}^{-3}$  at  $L = 2.5$ . We also find that for these conditions the  $\text{He}^+$  mass abundance is about 12–16% compared with  $\text{H}^+$ . Both decreases and increases in the  $\text{He}^+$  column abundance appear to be linearly correlated to changes in the percentage occurrence of  $\text{He}^+$  as determined from a combination of ground-based VLF and ULF observations, as well as in situ measurements from the RPI experiment.

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- G. J. Bailey and C. R. Wilford, Space and Atmosphere Research Group, Department of Applied Mathematics, University of Sheffield, Sheffield, S3 7RH, UK. (g.bailey@sheffield.ac.uk; c.wilford@sheffield.ac.uk)
- D. L. Carpenter, STAR Laboratory/Electrical Engineering Department, Stanford University, Stanford, CA 94305, USA. (dlc@nova.stanford.edu)
- M. A. Clilverd and M. C. Rose, British Antarctic Survey, Madingley Road, Cambridge, CB3 0ET, UK. (macl@bas.ac.uk; mcr@bas.ac.uk)
- J. Goldstein, Rice University, Rice Space Institute MS-108, 6100 Main Street, Houston, TX 77005-1892, USA. (jerru@hydra.rice.edu)
- I. R. Mann, Department of Physics, University of York, Heslington, York, YO10 5DD, UK. (ian@aurora.york.ac.uk)
- F. W. Menk, School of Mathematical and Physical Sciences and Cooperative Research Centre for Satellite Systems, University of Newcastle, Callaghan, N.S.W., 2308, Australia. (fred.menk@newcastle.edu.au)
- G. Milinevski, Ukrainian Antarctic Center, 16, blvd Tarasa Shevchenka, 01601, Kyiv Ukraine. (antarc@carrier.kiev.ua)
- B. W. Reinisch, Center for Atmospheric Research, University of Massachusetts, Lowell, 600 Suffolk Street, Lowell, MA 01854, USA. (bodo\_reinisch@uml.edu)
- B. R. Sandel, Lunar and Planetary Laboratory, University of Arizona, 1040 East Fourth Street, Room 901, Tucson, AZ 85721, USA. (sandel@arizona.edu)
- N. R. Thomson, Department of Physics, University of Otago, Dunedin, New Zealand. (thomson@physics.otago.ac.uk)
- K. H. Yearby, Space Instrumentation Group, Department of ACSE, University of Sheffield, Mappin Street, Sheffield, S1 3JD, UK. (k.yearby@sheffield.ac.uk)